

Are the electromagnetic constants really constant ?

Do they depend upon the vacuum conditions ?

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On behalf of the FLOWER collaboration

(F. Couchot, A. Djannati-Atai, M. Urban)

Workshop Quantum Vacuum and Gravitation

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What is the physical origin of the electromagnetic constants c , ϵ_0 and μ_0 ?

➤ The electrodynamical “constants” c , ϵ_0 and μ_0 are considered to be **fundamental constants**

⇒ There is no physical mechanism explaining their origin

⇒ They are assumed to be invariant in space and in time

➤ Part I

- ✓ We propose a mechanism where ε_0 , μ_0 and c originate from the properties of the quantum vacuum and its interaction with photons

Urban, Couchot, Sarazin, Djannati Atai, Eur. Phys. Journal D 67, 3 (2013) 58

- ✓ Two main consequences:

⇒ stochastic fluctuations of c are expected

⇒ ε_0 , μ_0 and c can vary if the conditions of the vacuum vary

➤ Part II

- ✓ Analogy of GR with a static space-time metric and variable c and m
- ✓ Application to cosmological redshift of supernova with a static (non expanding)

space-time metric

⇒ Possible interpretation of apparent Λ due to a time variation of the vacuum conditions

An effective description of the quantum vacuum

Vacuum filled with continuously appearing and disappearing **ephemeral** fermion pairs (f, \bar{f})

$$\text{Life time of the pair: } \tau = \frac{\hbar}{2} \times \frac{1}{\text{Energy borrowed from vacuum}}$$



An effective description of the quantum vacuum

Vacuum filled with continuously appearing and disappearing *ephemeral* fermion pairs (f, \bar{f})

➤ Average energy of a pair

$$W_f = K_W 2E_{rest} = K_W 2m_f c_{rel}^2$$

➤ Life time of the pair

$$\tau_f = \frac{\hbar}{2W_f} = \frac{1}{K_W} \frac{\hbar}{4m_f c_{rel}^2}$$

➤ Density of the pairs (quantum mechanic)

$$N_f \approx \frac{1}{\Delta x^3} \approx \left(\frac{2\pi\hbar}{\Delta p} \right)^3 \approx \left(\frac{\sqrt{K_W^2 - 1}}{\lambda_{C_f}} \right)^3$$

➤ The global electric charge, color and kinetic moment are null

But the electric and magnetic dipole moments are not null

➤ Only the charged fermions are considered (leptons & quarks) since we only study electromagnetic constants. However neutral fermions and bosons are also present.



K_W is the single free parameter in this model

Three distinct definitions for the speed of light in vacuum

➤ C_{rel} : maximal speed in special relativity $\Rightarrow E_{rest} = mc_{rel}^2$

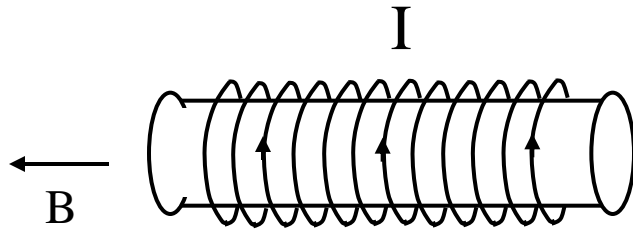
➤ $C_{E.M.}$: phase velocity of the E.M. wave $\Rightarrow c_{E.M.} = \frac{1}{\sqrt{\epsilon_0 \mu_0}}$

➤ C_γ : velocity of the photon $\Rightarrow c_\gamma = \frac{L_{propag}}{T_{propag}}$

A priori, we have on *average* : $c_{rel} = c_{E.M.} = c_\gamma$

μ_0

Vacuum Permeability μ_0



$$\mathbf{B} = \mu_0 \times (\mathbf{nI} + \mathbf{M})$$

\mathbf{M} = magnetization of matter

If the matter is removed : $\mathbf{B} = \mu_0 \mathbf{nI} \neq 0 !!!$

The vacuum is “globally” paramagnetic

In our model, μ_0 comes from the magnetization of the $f\bar{f}$ pairs

➤ We assume that the global kinetic moment of the pair is null

⇒ spins (fermion, antifermion) = $\uparrow\downarrow$ ou $\downarrow\uparrow$

➤ But opposite charges ⇒ the pair has a global magnetic moment = $2 \times$ Bohr magneton

$$2\mu_f = \frac{2eQ_f \hbar}{2m_f}$$

➤ When an external magnetic field B is applied:

⇒ pairs with magnetic moments aligned with B , have a longer life-time τ_f

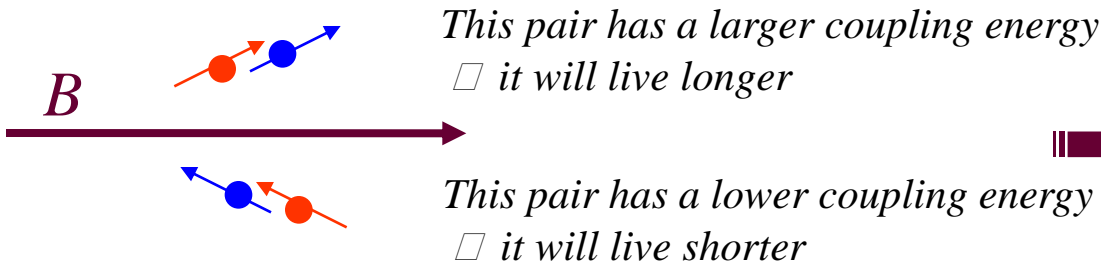
- The life-time τ_f of the pair depends on its coupling energy with B :

$$\tau_f(\theta) = \frac{\hbar/2}{W_f + W_{\text{coupling}}}$$

$$W_{\text{coupling}} = -2\mu_f B \cos\theta$$



$$\tau_f(\theta) = \frac{\hbar/2}{W_f - 2\mu_f B \cos\theta}$$



The difference of the life-times leads to a global magnetization of the vacuum

- By Averaging over θ $\implies \langle \mathcal{M}_i \rangle = \frac{\int_0^\pi 2\mu_i \cos\theta \tau_i(\theta) 2\pi \sin\theta d\theta}{\int_0^\pi \tau_i(\theta) 2\pi \sin\theta d\theta} \simeq \frac{4\mu_i^2}{3W_i} B.$

- It lead to a density of magnetization $M_i = 2N_i \langle \mathcal{M}_i \rangle$

- By summing over all the fermions (3 families)

$$\implies \frac{1}{\tilde{\mu}_0} = \sum_i \frac{M_f}{B} = c_{\text{rel}}^2 e^2 \sum \frac{2N_f Q_f^2 \lambda_{\text{Cf}}^2}{3W_f}$$

$$\left. \begin{aligned} W_f &= K_W 2m_f c_{rel}^2 \\ N_f &= \left(\frac{\sqrt{K_W^2 - 1}}{\lambda_{c_f}} \right)^3 \\ \lambda_f &= h / (m_f c_{rel}) \end{aligned} \right\} \Rightarrow \tilde{\mu}_0 = \frac{K_W}{\left(\sqrt{K_W^2 - 1} \right)^3} \times \frac{24\pi^3 \hbar}{c_{rel} e^2 \sum_f Q_f^2}$$

We must sum over the 3 charged leptons and the 6 quarks with 3 colors \Rightarrow
 $3+6 \times 3 = 21$ types de fermions.

$$\Rightarrow \sum_f Q_f^2 = e^2 \times \left(3 \times 1 + 3 \times 3 \times \left(\frac{4}{9} + \frac{1}{9} \right) \right) = 8e^2$$

$$\Rightarrow \tilde{\mu}_0 = \frac{K_W}{\left(\sqrt{K_W^2 - 1} \right)^3} \times \frac{3\pi^3 \hbar}{c_{rel} e^2}$$

$$\tilde{\mu}_0 = \mu_0 = 4\pi \cdot 10^{-7} \text{ N.A}^{-2} \Rightarrow \frac{\left(\sqrt{K_W^2 - 1} \right)^3}{K_W} = \frac{3\pi^2}{4\alpha} \Rightarrow \mathbf{K_W \approx 32}$$

The average energy of fermion pairs is ~ 32 times their mass energy ($2mc^2$)

Why $K_W \sim 32$?

If the energy spectrum density of the pairs $f\bar{f}$ is $p(E) = \frac{1}{E^2}$

$$\langle E_f \rangle = \frac{\int_{2mc^2}^{E_{Planck}} E \cdot p(E) \cdot dE}{\int_{2mc^2}^{E_{Planck}} p(E) \cdot dE} \approx \ln\left(\frac{E_{Planck}}{2m_f c^2}\right) \times 2m_f c^2$$



$$\begin{cases} K_W \sim 51 \text{ for } e^+e^- \\ K_W \sim 43 \text{ for } t\bar{t} \end{cases}$$

ε_0

Vacuum Permittivity ϵ_0

The mechanism is similar to $\mu_0 \Rightarrow \epsilon_0$ due to the polarization of the pairs $f\bar{f}$ in vacuum

➤ Electric dipole moment of the pairs $f\bar{f}$ $d_i = Q_i e \delta_i$ (δ_i is the average size of the pair)

➤ Pairs are polarized during their lifetime τ

➤ τ depends on the coupling energy of the pair with the electrostatic field E

$$\tau_i(\theta) = \frac{\hbar/2}{W_i - d_i E \cos \theta}$$

➤ τ is larger when the pair is aligned with $E \Rightarrow$ **POLARISATION**

$$\Rightarrow D = \tilde{\epsilon}_0 \times E \quad \text{with} \quad \tilde{\epsilon}_0 = e^2 \sum 2N_i Q_i^2 \frac{\delta_i^2}{3W_i} \quad K_W = 32 \quad \Rightarrow \tilde{\epsilon}_0 = \epsilon_0$$

➤ Finally, we can verify that

$$c_{E.M.} = \frac{1}{\sqrt{\tilde{\epsilon}_0 \tilde{\mu}_0}} = c_{rel}$$

**Let's see how a « real » photon would propagate
through this vacuum filled by « ephemeral » fermions**

Interaction of a photon with the fermion pairs in vacuum

- Real photon is trapped by an ephemeral pair
- As soon as the pair disappears, the photon is relaxed with its initial energy-momentum



- Between two interactions, the vacuum is « empty »
 - ⇒ there is no length scale, neither time scale
 - ⇒ the photon goes *instantaneously* to the next interaction
- The duration of the capture \approx the lifetime of the pair \Rightarrow finite transit time of the photon
 - ⇒ finite velocity
- A photon of a given helicity interact only with a fermion of opposite helicity (in order to flip its spin)

Derivation of the photon velocity c_γ

➤ σ_f = cross-section for a photon to interact with a $f\bar{f}$ pair

➤ When a photon crosses a length L of vacuum

The average number of stops on the $f\bar{f}$ pairs is $N_{stop,f} = L \times N_f \times \sigma_f$

And the average duration of stops on the $f\bar{f}$ pairs is $\bar{T}_f = N_{stop,f} \times \frac{\tau_f}{2}$

➤ The average total duration for a photon to cross a length L is $\bar{T} = \sum_f N_{stop,f} \times \frac{\tau_f}{2}$

➤ One obtains the general expression of the photon velocity in vacuum:

$$c_\gamma = \frac{L}{\bar{T}} = \frac{1}{\sum \sigma_f \times N_f \times \tau_f / 2}$$

Derivation of the photon velocity c_γ

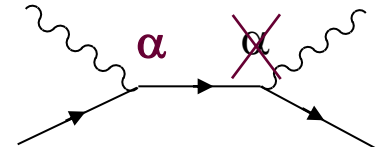
$$c_\gamma = \frac{L}{T} = \frac{1}{\sum \sigma_f \times N_f \times \tau_f / 2} \Rightarrow \left. \begin{array}{l} \tau_f = \frac{1}{K_W} \frac{\hbar}{4m_f c_{rel}^2} \\ N_f = \left(\frac{\sqrt{K_W^2 - 1}}{\lambda_{C_f}} \right)^3 \\ K_W \approx 32 \end{array} \right\} \Rightarrow c_\gamma = \frac{64 \alpha}{3\pi} \times \frac{1}{\sum \sigma_f^2 \lambda_{C_f}^2} \times c_{rel}$$

We can show that:

$$c_\gamma = c_{rel}$$

if

$$\sigma_f = 4 \times \frac{\sigma_{\text{Thomson}}}{\alpha}$$



We get a complete coherent model:

$$\langle c_\gamma \rangle = c_{E.M.} = c_{rel}$$

$$\left\{ \begin{array}{l} \langle c_\gamma \rangle = \text{average velocity of the photon} \\ c_{E.M.} = \frac{1}{\sqrt{\mu_0 \epsilon_0}} \\ c_{rel} : E = m c_{rel}^2 \end{array} \right.$$

Stochastic fluctuations of the speed of light

The successive interactions of the photon are independent

⇒ The number of captures and their duration fluctuate statistically

⇒ The propagation time of a photon to cross a length L of vacuum must fluctuate as

$$\sigma_t(L) = \sqrt{L} \times \frac{1}{c} \times \sqrt{\frac{\lambda_{ce}}{96\pi K_W}}$$

$$K_W \approx 32 \Rightarrow \sigma_t(L) = 50 \text{ as} \times \sqrt{L(\text{m})}$$

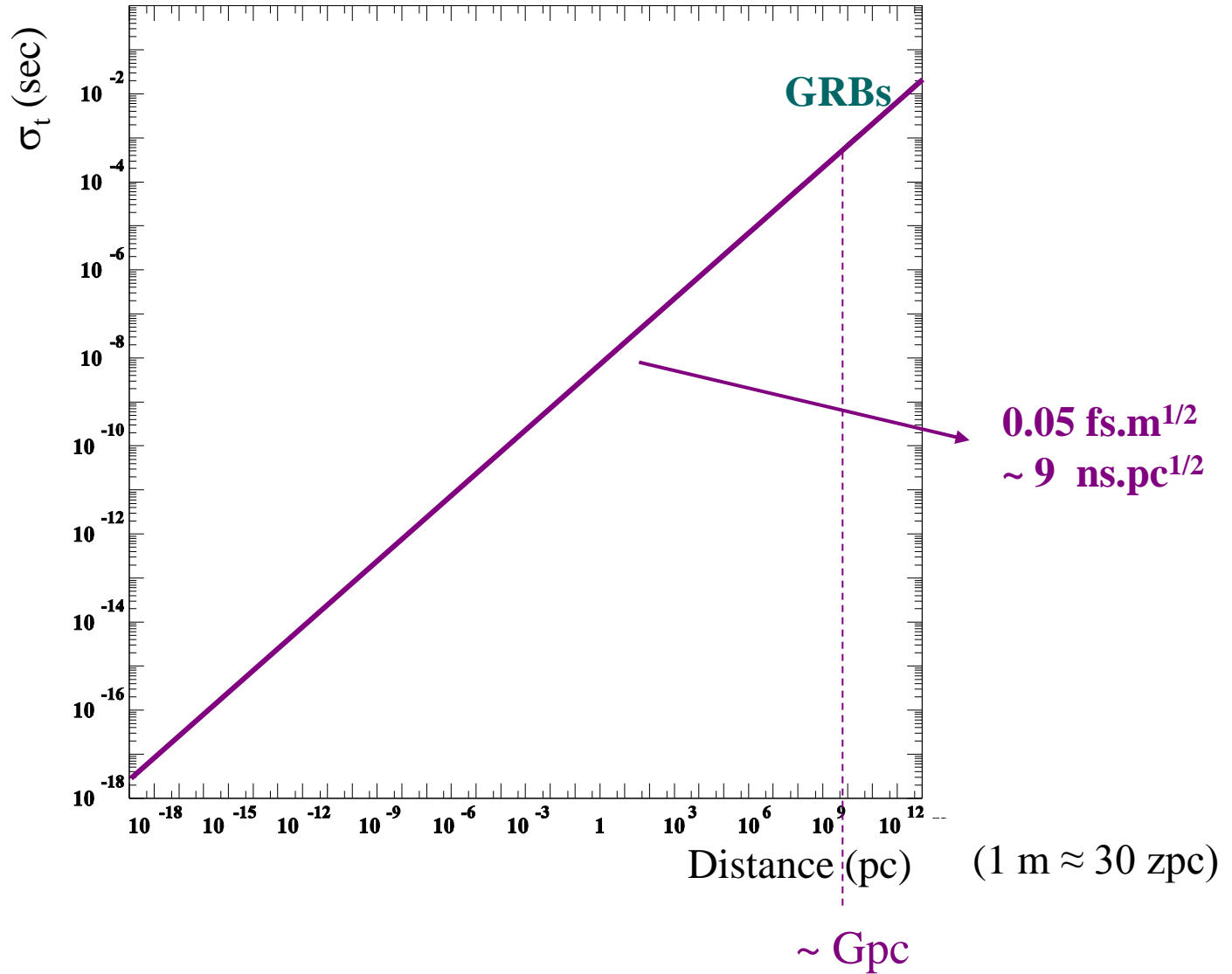
Search for a broadening of the time width of a light pulse
as the square root of the transit length

Remarks:

- No dispersion in frequency is expected (energy of the photon is conserved)
- Phase fluctuations are expected to be much lower (see “*reply to comment*” in

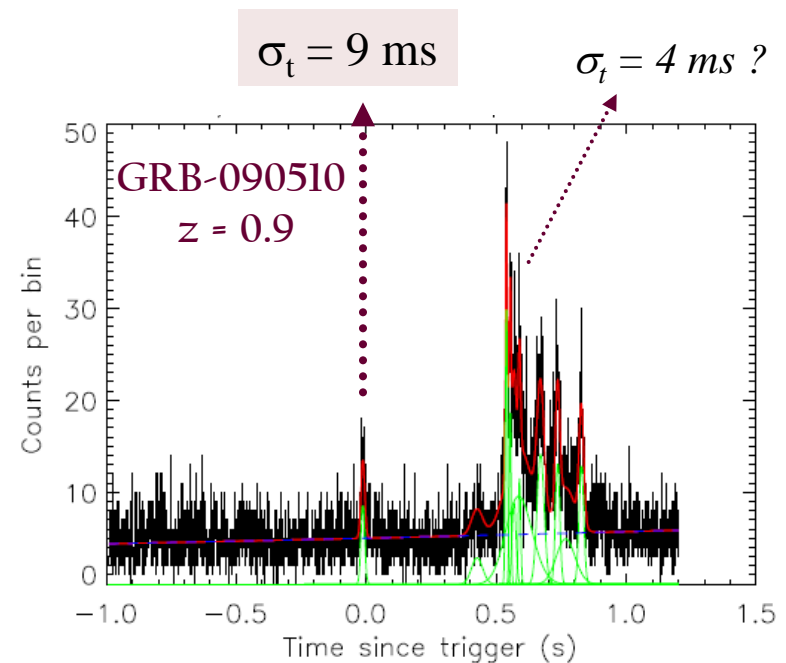
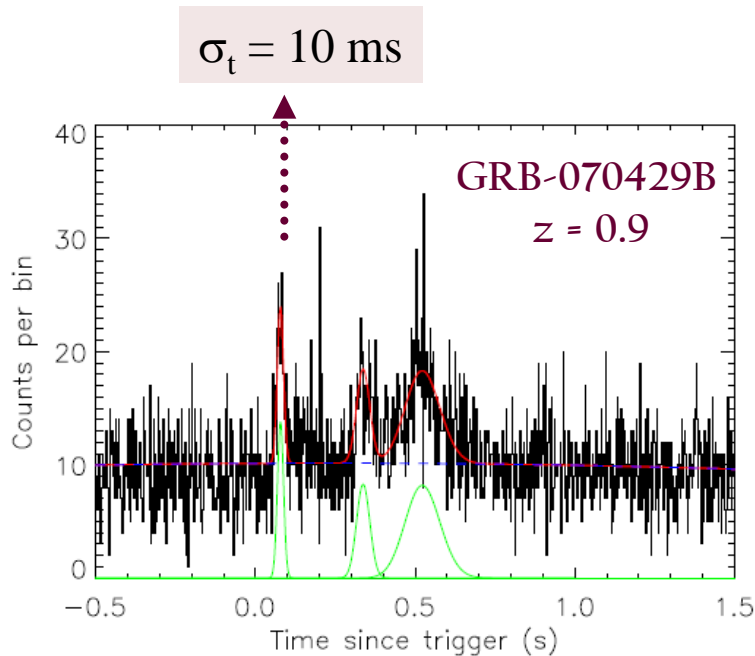
EPJ 2013)

Available constraints



Gamma Ray Bursts

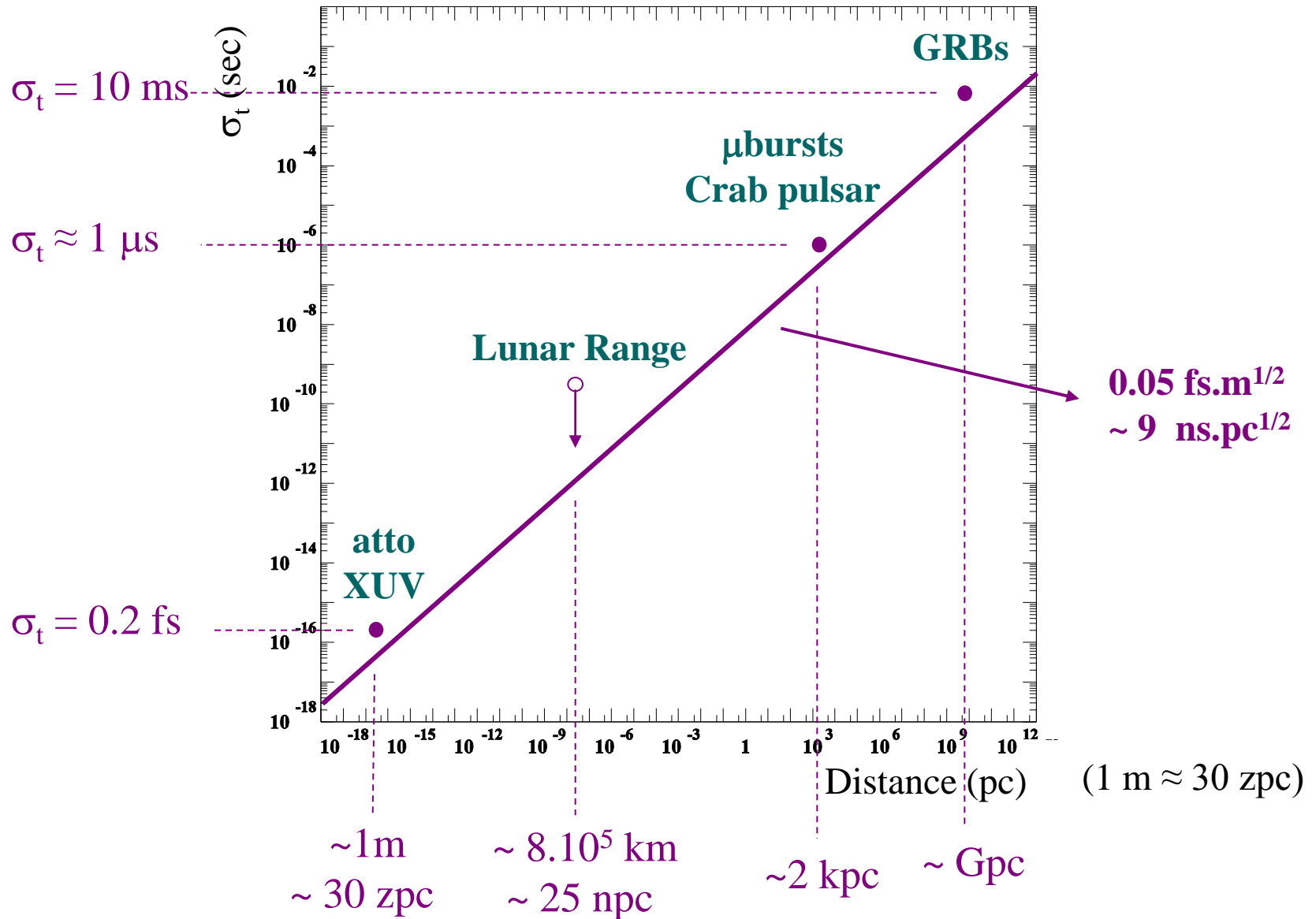
- ~ 20 short GRB's have been observed by SWIFT, Konus-Wind of FERMI with a reliable measured redshift
- An analysis of their light curve is in progress, in coll. with N. Bhat (Univ. Alabama in Huntsville)
- Preliminary results (after analysing 7 GRBs):



$$z = 0.9 \Rightarrow d_L \approx 2.10^{26} \text{ m}$$

$$\sigma_0 \leq 750 \text{ as.m}^{-1/2}$$

Available constraints



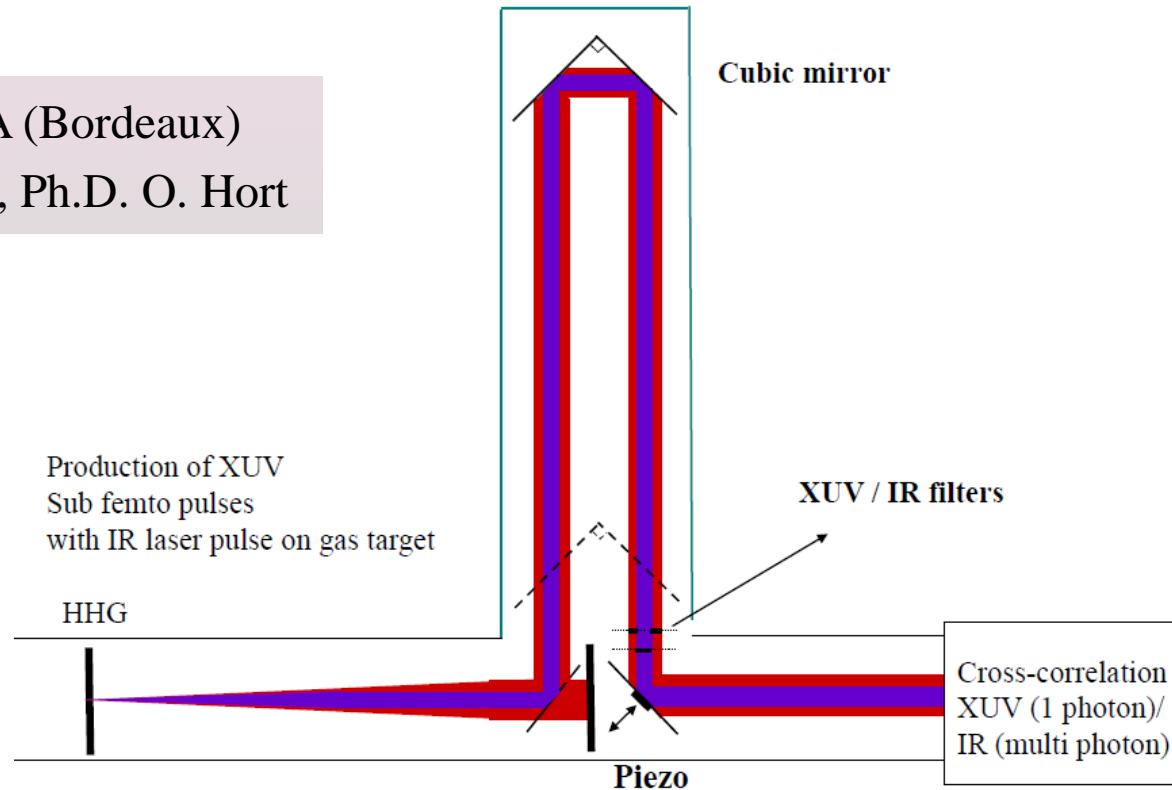
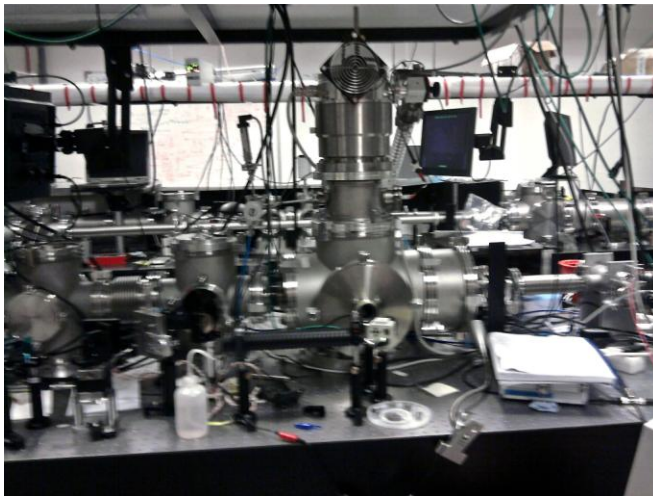
The *atto-FLOWER* experimental project

We propose to measure the duration of attosecond XUV pulses
after crossing few tens of meter of vacuum

$$\text{Assuming XUV pulse } \sigma_t = 0.1 \text{ fs} \xrightarrow[\sigma = 0.05 \text{ fs}\cdot\text{m}^{1/2}]{2 \times 25 \text{ m}} \sigma_t \approx 0.4 \text{ fs} ?$$

New collaboration with CELIA (Bordeaux)

E. Constant, E. Mevel, F. Catoire, Ph.D. O. Hort




PART II

Variation of the electromagnetic constants

- ϵ_0 , μ_0 and c are not constant but depend upon the vacuum parameters
⇒ they should vary when the vacuum parameters are modified by an external field
- They should vary when the vacuum is stressed by an external field
⇒ Gravitational field : analogy with the General Relativity
- They could vary in time if the vacuum also varies in time
⇒ Cosmology : analogy with the apparent expansion and its acceleration
- We will assume that e and \hbar are fundamental constants

$$\epsilon_0 = \text{cte} \times \frac{e^2}{\hbar c_{rel}} \Rightarrow \epsilon_0 \times c = \text{cte}$$


$$\alpha = \frac{e^2}{4\pi\epsilon_0\hbar c} = \text{cte}$$

The vacuum stressed by a gravitational field

- **General relativity:** geometrical theory based on the equivalent principle
 - ✓ The constants are constant
 - ✓ But we must modify the “mathematical” space-time coordinates

- **A possible physical analogy:**
 - ✓ We do not modify the “mathematical” space-time coordinates
 - ⇒ Static and flat space-time metric (no curvature)
 - ✓ But we accept that c and m are modified (the vacuum conditions are modified)
 - Variation of c : deflection of light, shapiro effect
 - Variation of m and c : gravitational redshift, perihelion shift

Deflection of the light

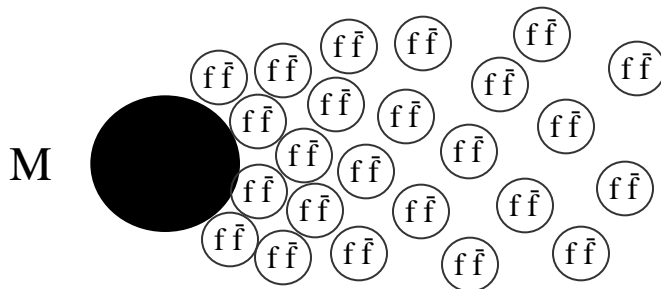
➤ **The analogy of GR with varying vacuum optical index in a flat space time is an old idea** to explain the deflection of light by a gravitational mass

- Eddington (*Space, Time and Gravitation*, p. 109 Cambridge University Press 1920)
- Landau & Lifshitz: (*The Classical Theory of Fields*, Pergamon Press 1975)
- Felice (*Gen. Rel. Grav.*, Vol. 2, 334 1971) & Evans, Nandi & Islam (*Gen. Rel. Grav.* 28, 413, 1996)

$$c(r) = \frac{c_\infty}{n(r)} \approx c_\infty \left(1 - \frac{2GM}{rc_\infty^2} \right)$$

c decreases when the photon get closer to the mass M

➤ **In this model, a mechanism exists to modify c**



$\left\{ \begin{array}{l} N_f \text{ is larger close to the mass } M \\ t_f \text{ is larger close to the mass } M \end{array} \right.$

$$\Rightarrow c(r) = c_\infty \times \frac{N_\infty \tau_\infty \sigma_\infty}{N(r) \tau(r) \sigma(r)} < c_\infty$$

Gravitational Blueshift

➤ Generalization for gravitational redshift, and other GR effects with variable optical index have been proposed

See: Evans, Nandi & Islam (*Gen. Rel. Grav.* 28, 413, 1996)


However:

- The energy of the photon was not conserved during its propagation
- The advance of the perihely was not correctly explained...

Gravitational Blueshift

In order to **conserve the energy of the photon during its propagation**, we also assume that the **inertial mass of a particle varies** when the vacuum is stressed by a gravitational field

$E_\gamma = pc = cte$
 c decreases
 p increases



$$E(r_1) = cte \times m_e(r_1) c^2(r_1)$$

$$E(r_2) = cte \times m_e(r_2) c^2(r_2)$$

Pound & Rebka :

$$E(r_2) = E(r_1) \times \left(1 - \frac{(r_1 - r_2) GM}{r_1 r_2 c_\infty^2} \right)$$



$$m(r) c^2(r) = m_\infty c_\infty^2 \left(1 - \frac{GM}{rc_\infty^2} \right) \Rightarrow m(r) = m_\infty \left(1 + \frac{3GM}{rc_\infty^2} \right)$$

Varying inertial mass

$$m(r) = m_{\infty} \left(1 + \frac{3GM}{rc_{\infty}^2} \right)$$

✓ This is what we expect if m corresponds to an electromagnetic self-energy with varying ϵ_0

(see Wilson *Phys. Rev.* 17, 54, 1921)

✓ Variation of the v.e.v. of the Higgs field ?

✓ Higgs doublet = condensate of $t\bar{t} + v_R v_L$ from vacuum fluctuation

\Rightarrow mass = self-energy ?

(see Smetana *arXiv:1309.4688*)

-

With varying c and m in a static and flat space-time metric

$$c(r) = \frac{c_\infty}{n(r)} = c_\infty \left(1 - \frac{2GM}{rc_\infty^2} \right)$$
$$m(r) = m_\infty \left(1 + \frac{3GM}{rc_\infty^2} \right)$$

we can fit well the observables

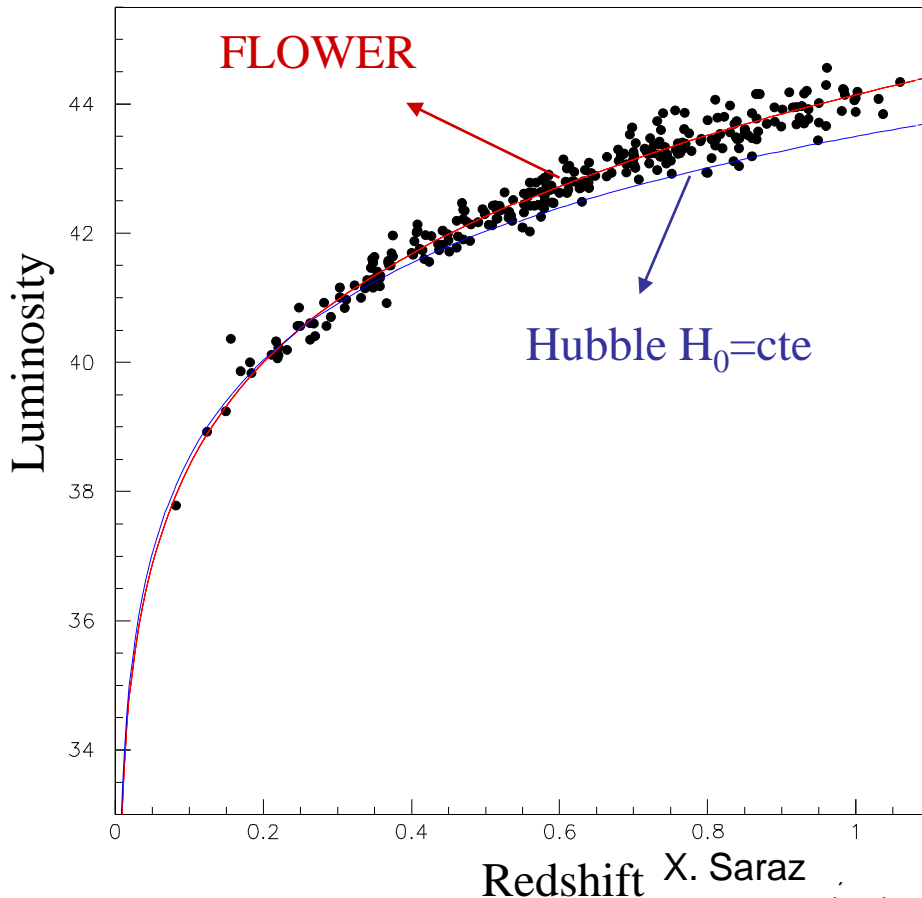
- Deflection of light
- Shapiro effect
- Gravitational redshift
- Perihelion shift of Mercury

Cosmological Redshift

We follow the same idea:

- ✓ Static & flat space-time metric (no expansion of the metric)
- ✓ Vacuum conditions modified in time: m and c vary in time

Supernovae Survey



$$c(t) = \frac{c_0}{n(t)} = \frac{c_0}{\sqrt{(1+t/\tau)}}$$

$$m(t) = m_0 (1+t/\tau)^{3/2}$$

t = epoch of the supernovae

Today $t = 0$

τ is the only free parameter of the fit

$\tau \sim 3.5 \text{ Gy}$

$$n(t) = \sqrt{(1+t/\tau)} \quad ?$$

What $n(t) = \sqrt{(1 + t / \tau)}$ could mean ?

“ $F=ma$ ” optics, Evans, *Am. J. Phys.* 54, 876 (1986)

Fermat's principle \equiv Principle of least action

$$\delta\left(\int n(r)dr\right)=0 \quad \equiv \quad \delta\left(\int L(r)dr\right)=0$$

$$\frac{n^2}{2} \quad \equiv \quad \phi$$

n = optical index

ϕ = potential energy

$$n(t) = \sqrt{(1 + t / \tau)} \quad \equiv \quad \phi(t) \propto 1 + \frac{t}{\tau}$$

The potential energy of the vacuum
decreases *linearly* in time

And do not ask me what it means !...

Conclusions

- ϵ_0 , μ_0 and c originate from the properties of the vacuum and its interaction with photons
- It is a discrete description of the vacuum
- Stochastic fluctuations of the photon propagation time in vacuum are predicted

$$\sigma_t(L) \approx 50 \text{ as} \times \sqrt{L(\text{m})}$$

We have started an experimental test in collaboration with CELIA using atto XUV pulses

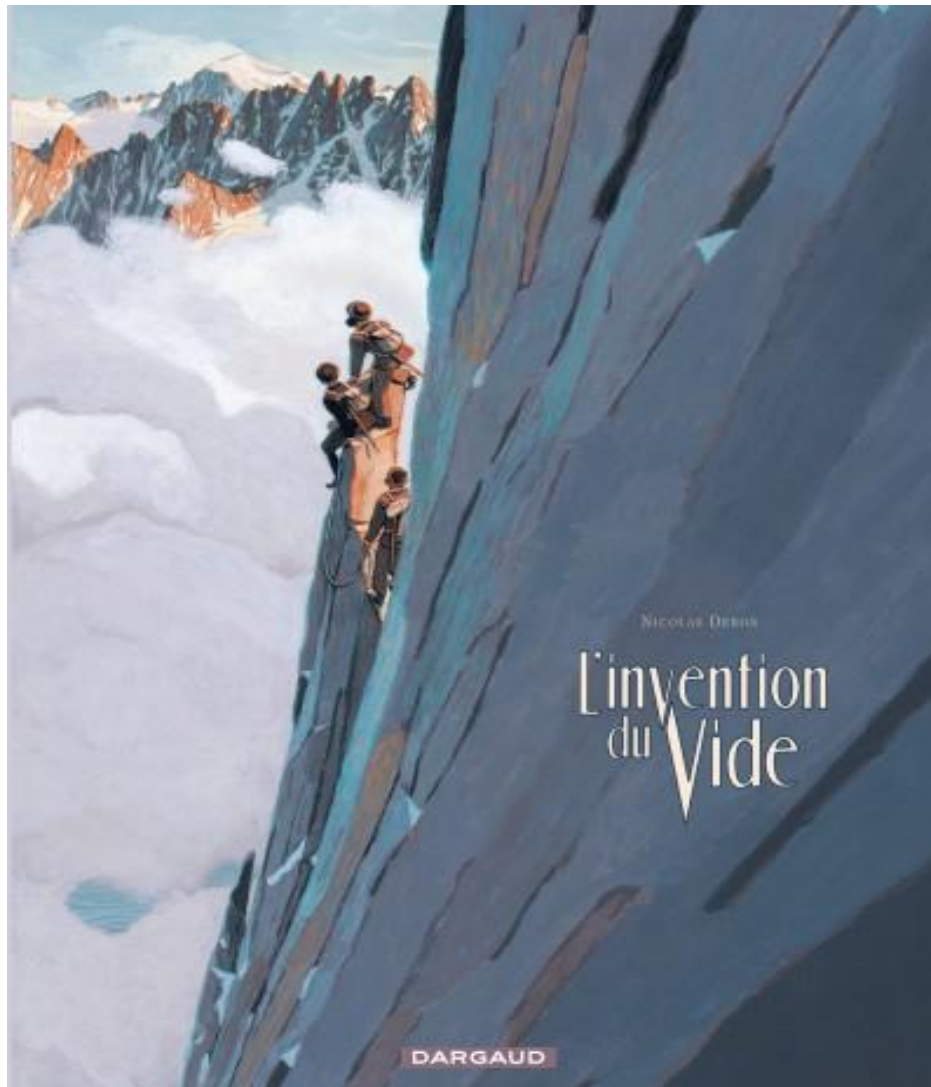
- ϵ_0 , μ_0 and c are not constant but can vary if the parameters of the vacuum vary
- Analogy of GR is possible with a static metric but with varying c and m
- Cosmology with a static space-time metric (no metric expansion) but time varying vacuum
 - ✓ Cosmological redshift is due to time variation of c and m
 - ✓ Apparent acceleration of Supernovae redshift is well fitted by a simple time variation of c equivalent to a linear vacuum potential variation

FLOWER

Fluctuations of the Light velocity Whatever the Reason

François Couchot, Arache Djannati-Atai, Xavier Sarazin, Marcel Urban

LAL Paris-Sud Orsay, APC Paris-Diderot

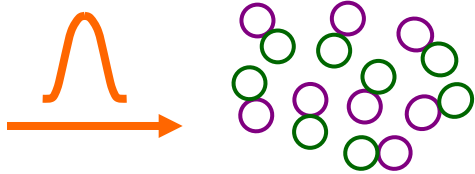


2nd experimental test

**The “vacuum optical index” n
 $n < 1$ inside an ultra high intensity laser pulse**

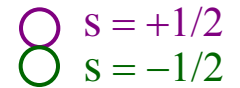
Effet SHADOK

The probe pulse, circularly polarized +1, moves with a velocity c

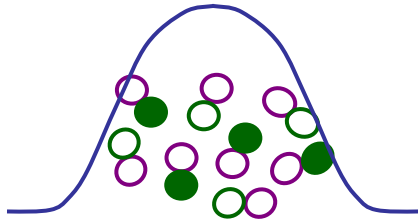


$$c = \frac{1}{\sum \sigma_i \times N_i \times \tau_i / 2} = 3.10^8 \text{ ms}^{-1}$$

Virtual fermion
antifermion pair

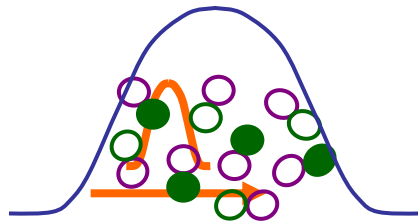


A pump laser, circularly polarized +1, with ultra high intensity, masks some virtual pairs



Number of occupied pairs $\Delta N = N_\gamma \sum (\sigma_i \times N_i)$

The probe pulse (circularly polarized +1) will move with a higher velocity c^*



$$c^* = \frac{1}{\sum \sigma_i \times (N_i - \Delta N_i) \times \tau_i / 2} \approx \frac{c}{1 - N_\gamma / (64N_e)}$$

$$\frac{\delta c}{c} = 1 - n \approx \frac{N_\gamma}{64N_e}$$

N_e is the e^+e^- pair density in vacuum

Light Saber

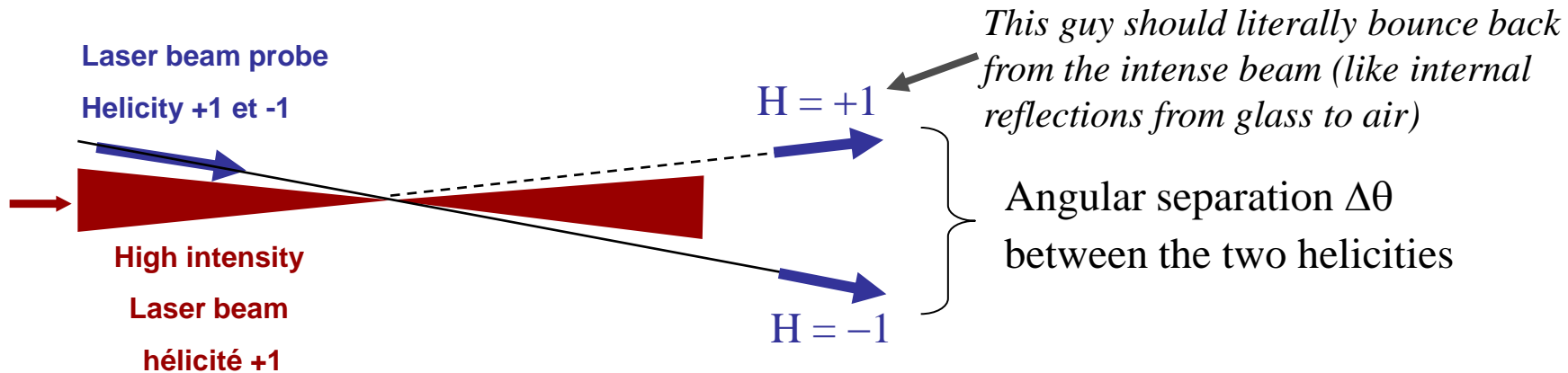
$$1 - n \approx \frac{N_\gamma}{64 N_e}$$

with $N_e \approx \left(\frac{\sqrt{K_W^2 - 1}}{\lambda_C} \right)^3 \approx \left(\frac{32}{\lambda_C} \right)^3 \approx \left(\frac{1}{76 \text{ fm}} \right)^3 \approx 2.10^{39} \text{ cm}^{-3}$

Laserix @ Orsay (40 J, 30 fs, 0.1 Hz)
Focalized in $10 \mu\text{m}^2$

$\Rightarrow \tilde{N}_\gamma \approx 10^{21} \text{ photons/pulse}$
 $V_{pulse} \approx 4 \cdot 10^{-17} \text{ m}^3$

$$1 - n \approx 2 \cdot 10^{-4}$$



$$\Delta\theta \approx \sqrt{2(1 - n)} \approx 1 \text{ mrad}$$